

I. Some mathematical preliminariesI.1 The Schrödinger equation

In quantum mechanics, the state of a particle is

described by a wave function $\psi: \mathbb{R}^3 \times \mathbb{R} \rightarrow \mathbb{C}$

with the interpretation that for fixed $t \in \mathbb{R}$

$|\psi(\cdot, t)|^2$ is the prob. distribution for the position of the particle.

The state space (for one particle) is the Hilbert-

space $L^2(\mathbb{R}^3) = \{ \psi: \mathbb{R}^3 \rightarrow \mathbb{C} / \int |\psi(x)|^2 dx < \infty \}$

with scalar product

$$\langle \psi, \phi \rangle = \int_{\mathbb{R}^3} \overline{\psi(x)} \phi(x) dx$$

The Schrödinger equation is given by

$$(1) \begin{cases} i\hbar \frac{\partial}{\partial t} \psi = H \psi \\ \psi(0) = \psi_0 \in L^2(\mathbb{R}^3) \\ \text{(i.e., } \psi(x, 0) = \psi_0(x) \text{)} \end{cases}$$

$$(2) \quad H \psi(x, t) = -\frac{\hbar^2}{2m} \Delta_x \psi(x, t) + V(x) \psi(x, t)$$

\nearrow Laplace $\Delta_x = \sum_{j=1}^3 \frac{\partial^2}{\partial x_j^2}$
 \nwarrow Potential.

Note that (1) yields an initial value problem, also called Cauchy problem.

$$\hbar = \text{Planck's constant} = \frac{h}{2\pi}$$

Domain $\mathcal{D}(H) \subset L^2(\mathbb{R}^3)$ is unspecified at the moment (it usually is the Sobolev space H^2 of functions having square integrable second derivatives)

Remark: From the interpretation of Q.M. we need

that

$$\langle \psi(t), \psi(t) \rangle = \int |\psi(x, t)|^2 dx = \int |\psi(x, 0)|^2 dx \quad \forall t \in \mathbb{R}$$

(conservation of probability).

This is the case if H is symmetric, i.e., if

for all $\psi, \phi \in \mathcal{D}(H)$ one has

$$\langle \psi, H\phi \rangle = \langle H\psi, \phi \rangle$$

Indeed, if $\psi(t)$ solves the Cauchy problem (1), with $\psi(0) = \psi_0 \in \mathcal{D}(H)$

then, with $\dot{\psi} = \frac{\partial}{\partial t} \psi$

$$\begin{aligned} \frac{d}{dt} \langle \psi(t), \psi(t) \rangle &= \langle \dot{\psi}, \psi \rangle + \langle \psi, \dot{\psi} \rangle \\ &= \left\langle \frac{1}{i\hbar} H \psi, \psi \right\rangle + \left\langle \psi, \frac{1}{i\hbar} H \psi \right\rangle \\ &= \frac{1}{i\hbar} \left(\langle \psi, H \psi \rangle - \langle H \psi, \psi \rangle \right) = 0 \end{aligned}$$

iff H is symmetric (and if $\psi(0) \in \mathcal{D}(H)$)

Note that we implicitly assumed here that H 's are nice:

we needed a) $\psi_0 \in \mathcal{D}(H) \Rightarrow \psi(t) \in \mathcal{D}(H)$

b) The Cauchy problem (1) has a solution \checkmark

In the following we consider the Cauchy-problem for an abstract linear operator H on a Hilbert-space \mathcal{H} .

Definition: We say the dynamics exists, if for all $\psi_0 \in \mathcal{H}$ the Cauchy problem (1) has a unique solution which conserves probability.

Theorem 2: The dynamics exists iff H is self-adjoint.

We want to prove this, but first we need to give meaning of most of the above words.

I.2 Some mathematical supplements

I.2.1 Spaces, norms, scalar products etc.

Vectors spaces: A vector space V is a collection of elements for which the operations of addition, $(u, v) \mapsto u+v$ and multiplication by a (real or complex) number, $\mathbb{F} \times V \rightarrow (a, v) \mapsto av$ $\mathbb{F} = \mathbb{R}$ or \mathbb{C} are defined such that

$$u+v = v+u \quad (\text{commutativity})$$

$$u+(v+w) = (u+v)+w \quad (\text{associativity})$$

$$u+0 = 0+u = u \quad (\text{existence of zero vector})$$

$$\alpha(\beta u) = (\alpha\beta)u$$

$$(\alpha+\beta)u = \alpha u + \beta u$$

$$\alpha(u+v) = \alpha u + \alpha v$$

$$0 \cdot v = 0, \quad 1 \cdot v = v$$

hold. We also denote $-v = (-1) \cdot v$

Elements of a vector space are called vectors

Examples

$$a) \mathbb{R}^d = \{x = (x_1, \dots, x_d) \mid x_j \in \mathbb{R} \quad j=1, \dots, d\}$$

Euclidean space of dimension d .

b) $\mathcal{C}(\Omega) =$ space of continuous functions on Ω
(real or complex valued)

c) $\mathcal{C}^k(\Omega) =$ space of k times cont. diff. fcts on Ω
($\Omega \subset \mathbb{R}^d$ open)

Norms: A norm on a vector space V is a map

$$V \ni u \mapsto \|u\| = \|u\|_V \in \mathbb{R}_+ = [0, \infty) \text{ s.t.}$$

$$a) \|\alpha u\| = |\alpha| \|u\|$$

$$b) \|u+v\| \leq \|u\| + \|v\| \quad (\Delta\text{-ineq.})$$

$$c) \|u\| = 0 \Leftrightarrow u = 0.$$

Ex: • $V = \mathbb{R}^d \quad \|x\| = |x| = \left(\sum_{j=1}^d x_j^2\right)^{1/2}$

• $\|f\|_\infty := \sup_{x \in \Omega} |f(x)|, \quad f \in \mathcal{C}(\Omega).$

• $\|f\|_{C^k} := \max_{0 \leq |\alpha| \leq k} \left\| \frac{d^\alpha}{dx^\alpha} f \right\|_\infty \quad \text{on } \mathcal{C}^k(\Omega)$

$$\bullet \|f\|_p = \|f\|_{L^p} := \left(\int_{\Omega} |f(x)|^p d\mu \right)^{1/p}$$

A normed vector space is a vector space equipped with a norm.

Examples • \mathbb{R}^d with $\|x\| = |x|$

• $\mathcal{C}_b(\Omega) \subset \mathcal{C}(\Omega)$, cont.-bounded fns on $\Omega \subset \mathbb{R}^d$
with norm $\|f\|_{\infty}$.

• $\mathcal{C}^k(\Omega)$ with $\|f\|_{\mathcal{C}^k}$ or $\|f\|_p$.

Banach spaces: A normed vector space is complete, if

Cauchy sequences converge, i.e., if $(u_n) \subset V$

with $\lim_{l,m \rightarrow \infty} \|u_l - u_m\| = 0$, then $\exists u \in V$ with

$$\lim_{l \rightarrow \infty} \|u - u_l\| = 0.$$

A complete normed vector space is called Banach space

Ex: • \mathbb{R}^d with norm $\|x\| = |x|$

• $\mathcal{C}_b(\Omega)$ with norm $\|f\|_{\infty}$

• $\mathcal{C}^k(\Omega)$ with norm $\|f\|_{\mathcal{C}^k}$.

• $1 \leq p < \infty$ $L^p(\Omega) := \left\{ f: \Omega \rightarrow \mathbb{C} \mid \int_{\Omega} |f(x)|^p dx < \infty \right\}$
with norm $\|f\|_p$

(also $L^\infty(\Omega) = \left\{ f: \Omega \rightarrow \mathbb{C} \mid \operatorname{ess\,sup}_{\Omega} |f(x)| < \infty \right\}$)

with

$$\|f\|_{\infty} := \operatorname{ess\,sup}_{x \in \Omega} |f(x)| =$$

$$= \inf \{ M \mid |f(x)| \leq M \text{ almost every } x \}.$$

Dual space

1) A bounded linear functional on a normed vector space $(V, \|\cdot\|)$ is a map $l: V \rightarrow \mathbb{C}$ (or $V \rightarrow \mathbb{R}$ if V is a real vector space) s.t. it is linear

$$l(\alpha u + \beta v) = \alpha l(u) + \beta l(v)$$

$\forall u, v \in V, \alpha, \beta \in \mathbb{C}$ (or \mathbb{R}) and there is

$$C < \infty \text{ s.t. } |l(u)| \leq C \|u\| \quad \forall u \in V.$$

2) $V^* =$ dual space of V

$:=$ space of all bounded lin. funct. on V .

Note that V^* has a canonical norm given by

$$\|l\|_{V^*} := \sup_{u \in V} |l(u)|$$

$$\|1\| = 1$$

and $(V^*, \|\cdot\|_{V^*})$ is always a Banach-space.

One often writes $\langle l, u \rangle := l(u)$.

Scalar products and Hilbert spaces

Let \mathcal{H} be a (complex) vector space, and assume that

\mathcal{H} has an inner-product (scalar product) $\langle \cdot, \cdot \rangle$

i.e., $\langle \cdot, \cdot \rangle : \mathcal{H} \times \mathcal{H} \rightarrow \mathbb{C}$ obeys

$$a) \langle v, \alpha w + \beta z \rangle = \alpha \langle v, w \rangle + \beta \langle v, z \rangle$$

linearity $\forall v, w, z \in \mathcal{H}, \alpha, \beta \in \mathbb{C}$.

$$b) \langle w, v \rangle = \overline{\langle v, w \rangle} \quad (\text{conjugate symmetry})$$

$$c) \text{ pos. definite } \langle v, v \rangle \geq 0 \text{ for all } v \in \mathcal{H}$$

$$\langle v, v \rangle = 0 \Rightarrow v = 0.$$

Note: • Cauchy Schwarz inequality: For $u, w \in \mathcal{H}$

$$|\langle u, w \rangle| \leq \|u\| \|w\|$$

$$\text{with } \|u\| := \sqrt{\langle u, u \rangle}.$$

• Because of Cauchy-Schwarz, $\|\cdot\|$ is a norm on \mathcal{H} .

If \mathcal{H} is complete under $\|\cdot\| := \sqrt{\langle \cdot, \cdot \rangle}$, then \mathcal{H} is called Hilbert space.

Ex: $L^2(\mathbb{R}^d) := \left\{ \psi: \mathbb{R}^d \rightarrow \mathbb{C} \mid \int_{\mathbb{R}^d} |\psi(x)|^2 dx < \infty \right\}$

$$\text{with } \langle \psi, \phi \rangle := \int_{\mathbb{R}^d} \overline{\psi} \phi dx = \int \overline{\psi} \phi$$

• Sobolev space of order n , $n \in \mathbb{N}$:

$$H^n(\mathbb{R}^d) = \left\{ \psi \in L^2(\mathbb{R}^d) \mid \partial^\alpha \psi \in L^2(\mathbb{R}^d) \forall \alpha, |\alpha| \leq n \right\}$$

$$\alpha = (\alpha_1, \dots, \alpha_d) \in \mathbb{N}_0^d \text{ multiindex}$$

$$|\alpha| := \sum_{j=1}^d \alpha_j$$

$$\partial^\alpha \psi = \partial_1^{\alpha_1} \partial_2^{\alpha_2} \dots \partial_d^{\alpha_d} \psi$$

with inner product

$$\langle \psi, \phi \rangle_{H^n} := \sum_{\alpha: |\alpha| \leq n} \langle \partial^\alpha \psi, \partial^\alpha \phi \rangle$$

With the help of the Fourier transform one can characterize the Sobolev space

$$\psi \in H^n(\mathbb{R}^d) \Leftrightarrow \int_{\mathbb{R}^d} (1+|x|^2)^n |\vec{\psi}(x)|^2 dx < \infty$$

$$\Leftrightarrow \int_{\mathbb{R}^d} (1+|x|^2)^{2n} |\vec{\psi}(x)|^2 dx < \infty,$$

Parseval relation

A set $\{v_n \mid n \in \mathbb{N}\} \subset \mathcal{H}$ is orthonormal if

$$\|v_n\| = 1 \text{ all } n, \quad \langle v_n, v_m \rangle = 0 \text{ if } n \neq m.$$

It is a complete orthonormal set (or a basis) if in addition

$$\overline{\text{span}(v_n \mid n \in \mathbb{N})} = \mathcal{H}$$

where $\overline{A} = \text{closure of } A \subset \mathcal{H}$.

= set of all limit points

$\text{span}(v_n \mid n \in \mathbb{N}) = \text{set of } \underline{\text{finite}} \text{ linear combinations}$
of the v_n .

Parseval relation: If $\{v_n\}_n$ is a complete orthonormal

set, then $\forall w \in \mathcal{H}$

$$\|w\|^2 = \sum_n |\langle w, v_n \rangle|^2$$

$$\begin{aligned}
 \text{For physics:} \\
 &= \sum_n \langle w, v_n \rangle \langle v_n, w \rangle \\
 &= \langle w | \sum_n |v_n\rangle \langle v_n| w \rangle
 \end{aligned}$$

$$\text{so } \sum_n |v_n\rangle \langle v_n| = \text{identity on } \mathcal{H}.$$

Riesz representation

If \mathcal{H} is a Hilbert space, then its dual \mathcal{H}^* can be identified with \mathcal{H} itself via the map

$$\mathcal{H} \ni u \mapsto l_u := \langle u, \cdot \rangle \in \mathcal{H}^*$$

$$l_u(v) = \langle u, v \rangle \quad \forall v \in \mathcal{H}.$$

Thm (Riesz) The above map is onto, i.e., for

any $l \in \mathcal{H}^*$ exists $u \in \mathcal{H}$ s.t. $l = l_u = \langle u, \cdot \rangle$.

I.2.2 Operators on Hilbert spaces

A linear operator (or skew operator) on a Hilbert space \mathcal{H} is a tuple $(A, D(A))$,

$D(A)$ = domain of $A \subset \mathcal{H}$ linear & dense subset.

and $A: \mathcal{D}(A) \rightarrow \mathcal{H}$ s.t. $\forall u, v \in \mathcal{D}(A), \alpha, \beta \in \mathbb{C}$

$$A(\alpha u + \beta v) = \alpha Au + \beta Av$$

Ex - identity $\mathbb{1} = \text{id}: \mathcal{H} \rightarrow \mathcal{H}$ $\mathcal{D}(\mathbb{1}) = \mathcal{H}$

• Multiplication by a coordinate

$$x_j: \mathcal{H} \rightarrow \mathcal{H}$$

$$\text{i.e., } (x_j \psi)(x) := x_j \psi(x).$$

$$\mathcal{D}(x_j) = \left\{ \psi \in L^2 \mid \int x_j^2 |\psi(x)|^2 dx < \infty \right\} \\ \subsetneq L^2(\mathbb{R}^d).$$

Note $\mathcal{C}_0^\infty(\mathbb{R}^d)$ = set of ∞ often diff. fns with compact support

is dense in $L^2(\mathbb{R}^d)$ and $\mathcal{C}_0^\infty(\mathbb{R}^d) \subset \mathcal{D}(x_j)$

so $\mathcal{D}(x_j)$ is dense in $L^2(\mathbb{R}^d)$.

• Momentum operator

$$p_j: \mathcal{H} \rightarrow \mathcal{H}$$

$$\mathcal{D}(p_j) = \left\{ \psi \in L^2 \mid \mathcal{D}_j \psi \in L^2 \right\}$$

$$= \left\{ \psi \in L^2 \mid \int \mathcal{D}_j^2 |\psi(\mathbf{q})|^2 d\mathbf{q} < \infty \right\}$$

again $\mathcal{C}_c^\infty(\mathbb{R}^d) \subset \mathcal{D}(p_\delta)$, so $\mathcal{D}(p_\delta)$ is dense
in L^2 .

• Laplace $\Delta: \psi \mapsto \sum_{j=1}^d \partial_j^2 \psi$

$$\begin{aligned} \mathcal{C}_c^\infty \subset \mathcal{D}(\Delta) &= \{ \psi \in L^2 \mid \Delta \psi \in L^2 \} \\ &= \{ \psi \in L^2 \mid \int_{\mathbb{R}^d} \partial_j^2 |\psi(x)|^2 dx < \infty \} \end{aligned}$$

• Multiplication by a function $V: \mathbb{R}^d \rightarrow \mathbb{R}$.

$$V: \psi \mapsto V\psi, \text{ i.e., } (V\psi)(x) = V(x)\psi(x).$$

$$\mathcal{D}(V) = \{ \psi \in L^2 \mid \int V(x)^2 |\psi(x)|^2 dx < \infty \}$$

$$\supset \mathcal{C}_c^\infty \text{ if } V \in L_{loc}^2$$

• Schrödinger operator

$$H: \psi \mapsto -\frac{\hbar^2}{2m} \Delta \psi + V\psi$$

$$\mathcal{D}(H) = \mathcal{D}(H) \cap \mathcal{D}(V).$$

• Integral operator

$$K: \psi \mapsto \int k(\cdot, y) \psi(y) dy$$

$$\text{i.e., } (K\psi)(x) := \int k(x, y) \psi(y) dy$$

$k: \mathbb{R}^d \times \mathbb{R}^d \rightarrow \mathbb{C}$ is called integral kernel of K

Def 3 An operator A on \mathcal{H} is bounded, if

$$\|A\| := \sup (\|A\varphi\| \mid \varphi \in D(A), \|\varphi\| = 1) < \infty$$

operator
norm

$$= \sup (\|A\varphi\| \mid \varphi \in D(A), \|\varphi\| \leq 1)$$

$$= \sup (\|A\varphi\| \mid \varphi \in D(A), \|\varphi\| < 1)$$

$\mathcal{L}(\mathcal{H}) (= \mathcal{B}(\mathcal{H})) =$ set of all bounded linear
operators.

Note $(\mathcal{L}(\mathcal{H}), \|\cdot\|)$ is a Banach space

Lemma 4 If an operator A satisfies $\|A\varphi\| \leq C\|\varphi\|$

for all φ in a dense domain $D \subset \mathcal{H}$, then it extends
uniquely to a bounded operator (also called A) on
all of \mathcal{H} with $\|A\varphi\| \leq C\|\varphi\| \quad \forall \varphi \in \mathcal{H}$.

idea of proof: If $u \in \mathcal{H} \Rightarrow \exists (u_n)_n \subset D, u_n \rightarrow u$

(D is dense in \mathcal{H}). So, since

$$\|Au_n - Au_m\| \leq C\|u_n - u_m\| \rightarrow 0 \quad n, m \rightarrow \infty$$

$(Au_n)_n$ is Cauchy in \mathcal{H} . So

$$v := \lim_{n \rightarrow \infty} Au_n \text{ exists}$$

Define $Au := v = \lim_{n \rightarrow \infty} Au_n$. Done.

Remark: One needs to check consistency, i.e., if

$u_n \rightarrow u$ and $\tilde{u}_n \rightarrow u$, then $v = \lim Au_n = \lim A\tilde{u}_n$

and that A is bounded with $\|A\| \leq C$.